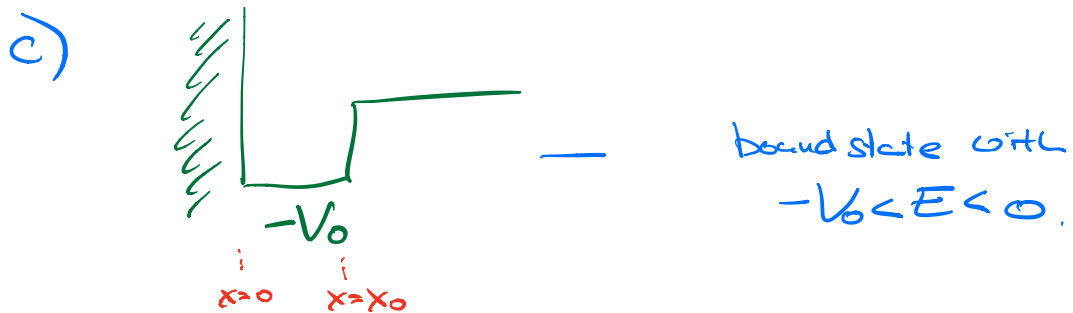


QP 1 - Oct. 29, 2020.

1) a) stat. state:  $H\psi = E \cdot \psi$   
 $H = \frac{p^2}{2m} \Rightarrow \psi = e^{ikx}$  (5pt)

b) linear superposition:  
 $\psi(x,t) = \frac{1}{\sqrt{2\pi}} \cdot \int dk \cdot \phi(k) \cdot e^{ikx} \cdot e^{-iEt/\hbar}$   
 (coeff. of linear superposition. (4pt))  
 $E = \frac{\hbar^2 k^2}{2m}$  (3pt)

normalisation:  
 $\int \psi^* \psi dx = \int \phi^* \phi dk = 1$  (3pt)



3)  $\psi = 0$  for  $x > x_0$   
 $\psi = C \cdot e^{-kx} + D \cdot e^{+kx}$  (1)  
 $E = -\frac{\hbar^2 k^2}{2m}$  (1)  
 $\psi = A \cdot e^{ikx} + B \cdot e^{-ikx}$  (1)  
 $E - V_0 = \frac{\hbar^2 k^2}{2m}$  (1)

d) Normalisation:  $D=0$

$$\Rightarrow \psi = C \cdot e^{-lx} \quad (2 \text{ pt})$$

$\psi$  Continuous @  $x=0$

$$A = -B \Rightarrow \psi = 2i \cdot A \cdot \sin(kx) \quad (4 \text{ pt})$$

$\psi$  and  $\psi'$  Continuous @  $x=0$ .

$$\begin{aligned} 2iA \cdot \sin(kx_0) &= C \cdot e^{-lx_0} & (2) \\ 2ik \cdot A \cdot \cos(kx_0) &= -l \cdot C \cdot e^{-lx_0} & (2) \end{aligned}$$

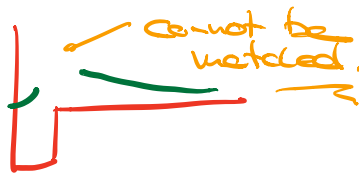
e) Small  $x_0 =$  Taylor expansion:

$$2iA \cdot k \cdot x_0 = C \quad (2)$$

$$2i \cdot k \cdot A = -l \cdot C \quad (2)$$

$$\Rightarrow x_0 = \frac{-1}{l} \quad (2)$$

so no bound state as  $x_0, l$  are positive. (4 pts)



2) a) Herm. operator:  $\langle A | O B \rangle = \langle O A | B \rangle$  (2pts)

take  $A=B$  eigenvector with eigenvalue  $a$ :

$$\begin{aligned} \langle A | O A \rangle &= a \cdot \langle A | A \rangle \\ &= \langle O A | A \rangle = a^* \cdot \dots \quad (3pts) \\ &\Rightarrow a = a^* \end{aligned}$$

b) Reality: (3pts)

\* guaranteed for  $\infty$ -dim VS & Herm. oper. with discrete spectrum

\* has to be assumed for -- (2pts)  
--- continuous ---

c) yes, entangled: (3pts)

measuring left particle influences Right Result. (2pts)

3) a)  $L_{\pm} = L_x \pm iL_y$      $[L_x, L_y] = i\hbar L_z$

$$[L_{\pm}, L_z] = [L_x \pm iL_y, L_z]$$

$$= [L_x, L_z] \pm i[L_y, L_z]$$

$$= -i\hbar L_y \pm i \cdot i\hbar L_x$$

$$= \mp i\hbar L_{\pm} \quad (5pts)$$

$$[L_{\pm}, L^2] = [L_x \pm iL_y, L_x^2 + L_y^2 + L_z^2]$$

$$= [L_x, L_y^2] + [L_x, L_z^2] \pm i \cdot [L_y, L_x^2] \pm i [L_y, L_z^2]$$

use  $[L_x, L_y^2] = L_x L_y^2 - L_y^2 L_x$   
 $= [L_x, L_y] L_y - L_y [L_y, L_x]$   
 $= i\hbar L_z L_y + i\hbar L_y L_z$

$$[L_x, L_z^2] = L_x L_z^2 - L_z^2 L_x$$

$$= [L_x, L_z] L_z - L_z [L_z, L_x]$$

$$= -i\hbar L_y L_z - i\hbar L_z L_y$$

cancel against each other.

$$\Rightarrow [L, L^2] = 0 \text{ and}$$

hence also  $[L_{\pm}, L^2] = 0$ .  
 (5pts)

b) if  $\psi$  is eigenstate of  $L^2, L_z$   
 with eigenvalues  $\lambda, \mu$ ,  
 then  $L_{\pm} \psi$  is also an eigenstate <sup>3pts</sup>  
 with eigenvalues  $\lambda, \mu \pm \hbar$ . <sup>2pts</sup>

c)  $l=1, m=0$ :  $\psi_0^0 = A \cos \theta$ .

act with  $L_{\pm}$ : (2pts)

$$\psi_0^{\pm} \sim \pm \hbar \cdot e^{\pm i\phi} \cdot \frac{\partial}{\partial \theta} (A \cos \theta)$$

$$\sim e^{\pm i\phi} \cdot \sin \theta$$

(3pts)

↑  
 ignoring  
 normalisation.

$$d) \quad L_+ L_+ \text{ on } Y_{1,0} \Rightarrow 0. \quad (5 \text{ pts})$$

explicitly:

$$\begin{aligned} & \left( \frac{\partial}{\partial \theta} + \frac{i}{\tan \theta} \frac{\partial}{\partial \phi} \right) e^{+i\phi} \sin \theta \\ &= e^{+i\phi} \cos \theta + i \cos \theta \cdot -i \cdot e^{+i\phi} \\ &= 0. \quad \checkmark \quad (5 \text{ pts}) \end{aligned}$$

$$e) \quad l_1 = 1 \quad \text{plus} \quad l_2 = 1$$

$\Rightarrow$  total angular momentum ranges from

$$\begin{array}{ccc} l_1 + l_2 & \dots & l_1 - l_2 \\ = 2 & & = 0 \end{array}$$

$$\Rightarrow l = 0, 1, 2. \quad (5 \text{ pts})$$

(or you can say that  $L^2 = l(l+1)\hbar^2 = 0, 2, 6 \cdot \hbar^2$  for full ang. mom.)